



# A Mohr circle method for 3D strain measurement using the geometry of no finite longitudinal strain and the $R_{XZ}$ strain ratio

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## ABSTRACT

This paper proposes a new method for 3D finite strain analysis. This method utilizes a Mohr circle construction combined with stereographic projection of the geometry of no finite longitudinal strain and with the strain ratio on the XZ-plane of the finite strain ellipsoid. The method is described using numerical examples and then it is tested by applying it to the deformed Deh Vazir conglomerate in the southwestern part of the Sanandaj-Sirjan HP-LT metamorphic belt, within the Zagros orogenic belt in Iran. The results of this method compare well with previous finite strain measurements using strain ratios on three principal planes of finite strain. Calculation of finite strain from strain ratios on the XY and YZ principal planes is advantageous when preparation of 3 perpendicular sections is difficult or impossible.

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## 1. Introduction

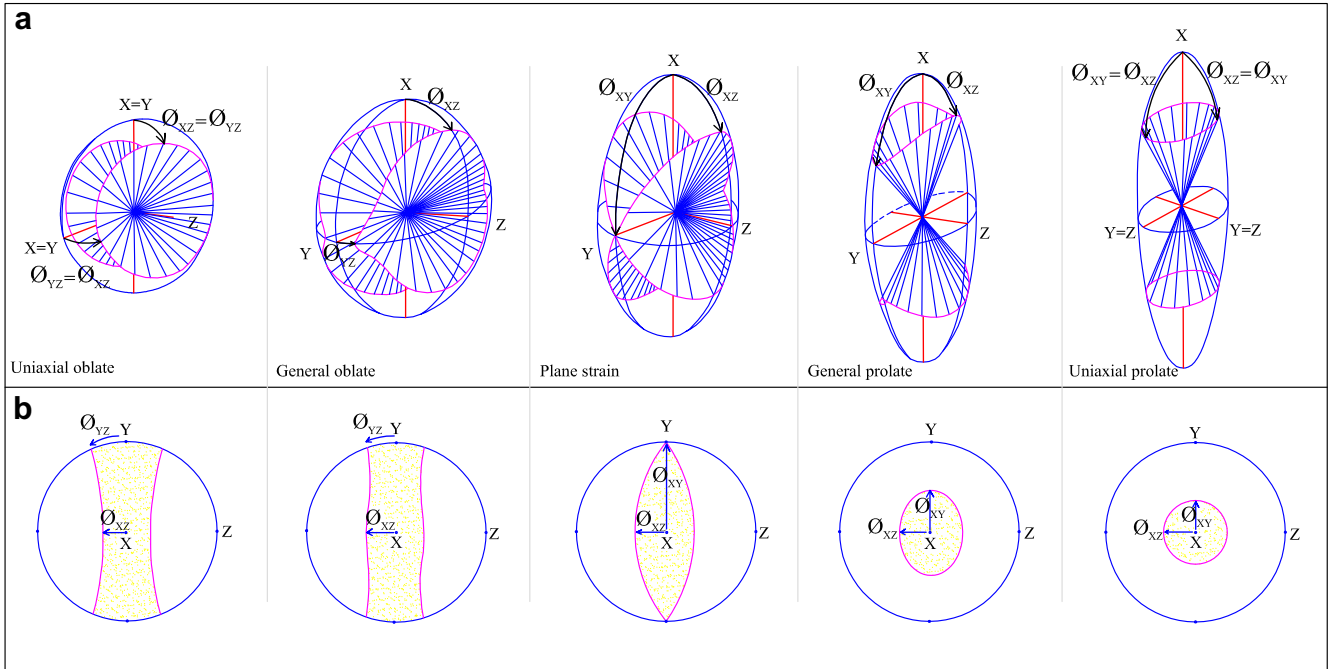
Measurements of finite strain and determination of the strain ellipsoid in naturally deformed rocks are important tasks for many structural geologists. The key to strain analysis lies in finding objects with known initial packing arrangement or features which enable final lengths or angles to be calculated. Following the classic paper of Cloos (1947), a diversity of methodologies has been proposed in order to estimate finite strain in deformed rocks. The  $R_f/\Phi$  method (Ramsay, 1967; Dunnet, 1969) and Fry method (Fry, 1979) are the most common methods that have been used by structural geologists; they use the shape ( $R_f/\Phi$  method) and distribution (Fry method) of objects (e.g. deformed ooids, pebbles of deformed conglomerate and deformed fossils) or of points (e.g. quartz grain centers in quartzite). Two dimensional finite strain can be completely described by three numbers which represent the orientations and magnitudes of the strain ellipse. Strain magnitude can be expressed as two principal stretch values, or a strain ratio and area change. For a full description of homogeneous 3D finite strain, six numbers are needed: three to describe the orientation of the strain ellipsoid, and three to describe strain magnitude. Strain magnitude can be expressed by the three principal stretch values, or by two strain ratios and volume change. Ramsay and Huber (1983, p. 198) suggested four approaches that can

be used to determine the 3D geometry of the strain ellipsoid from 2D strain ellipse data, among which the approach using two strain ellipses parallel to the principal planes is the most practical. This paper proposes a new method combining stereographic projection of orientations of no finite longitudinal strain together with the principal strain ratio on the XZ-plane and using a Mohr circle construction; the method allows one to determine  $R_{XY}$ ,  $R_{YZ}$  (strain ratio in the XY and YZ principal planes of the strain ellipsoid) and the geometry of the 3D strain ellipsoid. This method is applied to a deformed conglomerate to show how the finite strain varies on a variety of scales across a deformed area.

## 2. Geometry of no finite longitudinal strain

There are five main types of strain ellipsoid which can result from homogeneous deformation with no volume change (Flinn, 1962). Within the strain ellipsoid, lines whose deformed lengths are equal to their undeformed lengths are defined as lines of no finite longitudinal strain (n.f.l.s) (Ramsay, 1967). Orientations of all lines of no finite longitudinal strain define a surface separating sectors of positive and negative longitudinal strain within the ellipsoid. Except for the plane strain state where the strain ellipsoid shape ( $k$ ) = 1 which produces two circular sections, this surface is a double cone (on a circular or elliptical base) with common apices at the center of the ellipsoid. The shape of these conical surfaces can be completely described by the angles made by their lines of intersection on the principal planes and one of the principal strain axes ( $\Phi_{XY}$ ,  $\Phi_{XZ}$ ,  $\Phi_{YZ}$ ) (Fig. 1a). The values of

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**Fig. 1.** (a) Geometry of surfaces ( $k \neq 1$ ) and plane ( $k = 1$ ) of no finite longitudinal strain (n.f.l.s) in the five main types of the strain ellipsoid (Ramsay, 1967). (b) Half the maximum and minimum angular dimensions of either the extension field or shortening field are measured on a stereographic projection.

these angles depend upon the principal extensions. Half the maximum and minimum angular dimensions of either the extension field or shortening field are measured on a stereographic projection (Fig. 1b). For an oblate ellipsoid, these are measured from X in the XZ-plane ( $\Phi_{XZ}$ ) and from Y in the YZ-plane ( $\Phi_{YZ}$ ) so that  $\Phi_{XZ}$  is always greater than  $\Phi_{YZ}$ , unless the oblate ellipsoid is uniaxial ( $k = 1$ ), when  $\Phi_{XZ} = \Phi_{YZ}$ . For a prolate ellipsoid, the angular dimensions of the extension field are measured from X in the XZ-plane and designated  $\Phi_{XZ}$ , and from X in the XY-plane and designated  $\Phi_{XY}$ . These angles are equal for a uniaxial prolate ellipsoid, whereas  $\Phi_{XY} > \Phi_{XZ}$  for a triaxial prolate ellipsoid. The surface of no finite longitudinal strain is dependent on volume change; the two limits occur when the original sphere either lies completely outside or completely within the strain ellipsoid. Combining a study of deformed structures which have suffered elongation or compression (such as foliation, stretching lineation, boudins and folds) with stereographic analysis of the extension and compression directions will enable us to determine the pattern of no finite longitudinal strain. Comparing results of stereographic analysis with the no finite longitudinal strain pattern on a Flinn diagram (Flinn, 1962; Ramsay and Huber, 1983) provides a qualitative test for determining the homogeneity of strain. If the compression and extension domains are separated by irregular boundaries, the strain has been inhomogeneous. In such cases it is often possible to use data from smaller or larger domains to establish empirically the scale of any domains of homogeneous strain. Large bodies of inhomogeneously strained rock may be divided into volumetric units which are small enough to be considered to have deformed homogeneously (Talbot, 1970; Twiss and Moores, 1992). Similarly small domains of homogeneous strains may integrate within larger domains of homogeneous strains (Talbot, 1987; Talbot and Sokoutis, 1995).

2.1. The Mohr diagram for strain

Nadai (1950) first applied a Mohr circle to represent strain. Nadai recognized that a graph of  $\lambda'$  (reciprocal quadratic elongation) vs  $\gamma'$  ( $\gamma' = \gamma/\lambda$  or shear strain/quadratic elongation), in terms

of angles in the strain state, is identical in form with the Mohr stress diagram: the three principal planes of the strain ellipsoid are represented by three Mohr circles. The strain state for any other direction falls in the region bounded by the three circles. It was probably Brace (1961) who first used the expression “Mohr diagram” in the geological literature for the representation of three-dimensional finite strain. Ramsay (1967) expanded the usage of the Mohr diagram for strain in two ways. First, he illustrated the potential of Mohr circles for representing strain data. Second, he illustrated how the diagram for three-dimensional reciprocal strain could be used to calculate strain values. For five ellipsoid examples, Ramsay (1967) derived strain contours which were plotted on stereographic projections. In practice, it is more convenient to represent three-dimensional strain on a half Mohr diagram with three semicircles, thus representing the values, but not the sign, of  $\gamma'$ . Usually determination of the strain ellipsoid shape is possible by using the popular K factor (Flinn, 1962), where:

$$K = \left[ (\lambda_1/\lambda_2)^{1/2} - 1 \right] / \left[ (\lambda_2/\lambda_3)^{1/2} - 1 \right]$$

$K = 1$  (plane strain ellipsoid, type 3);  $K = 0$  (uniaxial oblate ellipsoid, type 1);  $K = \infty$  (uniaxial prolate ellipsoid, type 5);  $0 < K < 1$  (three axial oblate ellipsoid, type 2) and  $1 < K < \infty$  (three axial prolate ellipsoid, type 4). The Mohr diagrams are also a useful means of illustrating and classifying strain ellipsoids of different types. The relationship of the  $\psi_{max}$  (maximum shear angle) tangents to the three circles allows ellipsoids to be assigned to one of Ramsay’s five equal volume ellipsoid types (Fig. 2). The plane strain ellipsoids are immediately distinguishable with their  $\lambda'_1 \lambda'_2$  and  $\lambda'_2 \lambda'_3$  circles which share the same  $\psi_{max}$  line (Fig. 2c). The two end-member ellipsoids, uniaxial prolate ( $K = \infty$ ) and uniaxial oblate ( $K = 0$ ), are represented on the Mohr diagram by a single circle (Fig. 2a and e). In each case, one principal circle reduces to a point and the  $\psi_{max}$  is only defined in the XZ principal plane of the strain ellipsoid. In the three axial oblate and prolate ellipsoids there are different  $\psi_{max}$  for all principal planes of the strain ellipsoid (Fig. 2b and d).

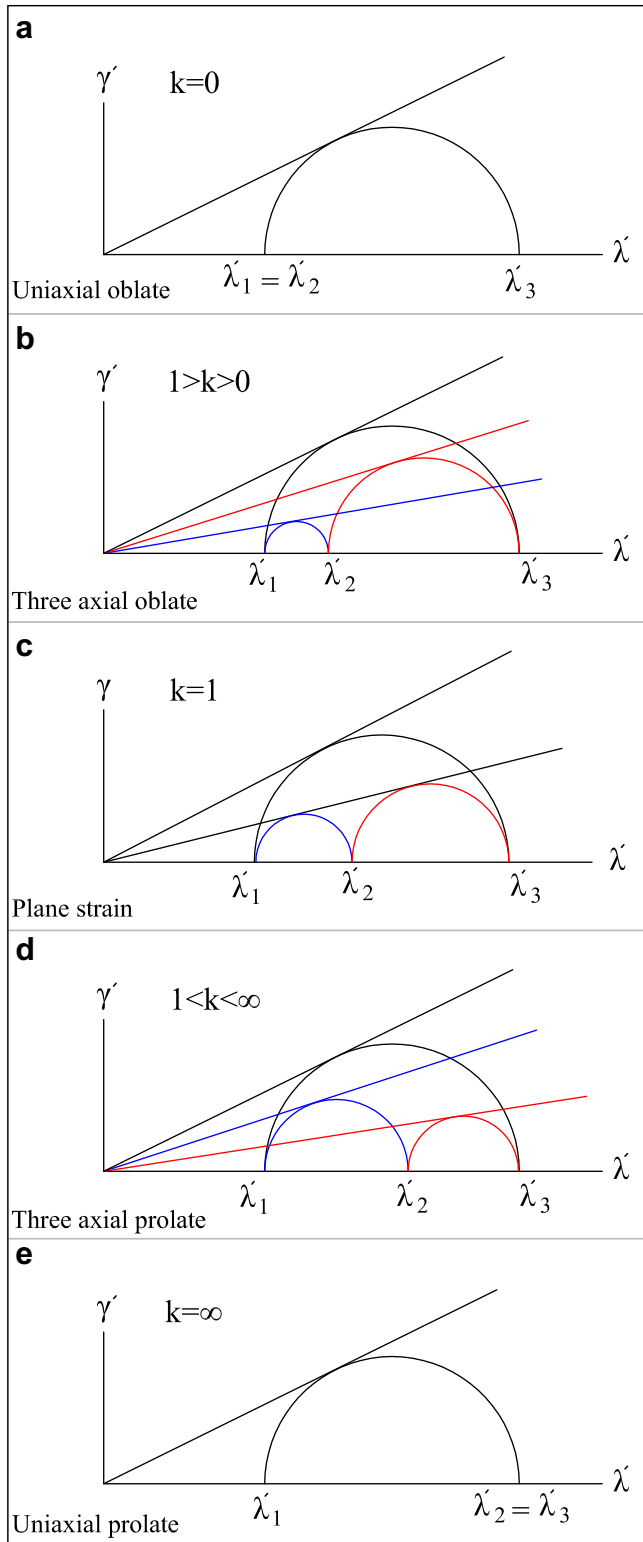


Fig. 2. The relationship of the  $\psi_{\max}$  tangents to the three principal circles allows strain ellipsoids to be assigned to one of Ramsay's five equal volume ellipsoid types.

## 2.2. Numeric example

Let us suppose that the data in Table 1 show the strain ratio in the XZ-plane of the strain ellipsoid and the  $\Phi$  angles obtained from strain analysis and structural studies in three geological settings. It is

**Table 1**

Numerical data for three cases showing the strain ratio in the XZ-plane of the strain ellipsoid and  $\Phi_{XY}$ ,  $\Phi_{XZ}$ ,  $\Phi_{YZ}$  angles. Incomplete data are assumed: in each case it is assumed that one of the angles,  $\Phi_{XY}$ ,  $\Phi_{XZ}$  and  $\Phi_{YZ}$ , is missing.

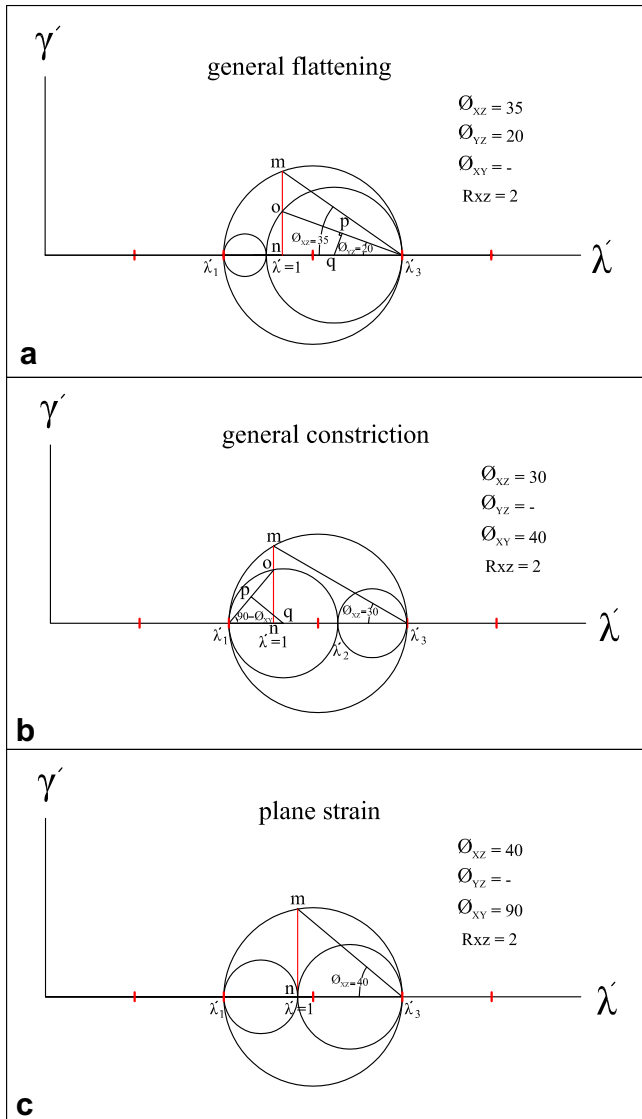
	$R_{XZ}$	$\Phi_{XZ}$	$\Phi_{XY}$	$\Phi_{YZ}$
Case 1	2	35	–	20
Case 2	2	30	40	–
Case 3	2	40	90	–

possible to determine the strain ratio in the other two principal planes ( $R_{XY}$  and  $R_{YZ}$ ) by using a Mohr circle construction. The data of Table 1 were used to calculate the strain ratio in the XY- and YZ-planes for states 1 to 3, respectively. In the first case the values of  $\Phi$  indicate an oblate geometry for the strain ellipsoid. To measure the strain ratio in the other two principal planes based on a Mohr circle construction, the following steps are used. Draw perpendicular axes  $c\lambda'$  and  $c\gamma'$ , where  $c$  is a constant of unknown value (Fig. 3a). For strain analysis of naturally deformed rocks, because it is rare to know the original size of the strain markers, we can only determine the relative strain ratios of the principal planes of finite strain ellipsoid but not the absolute strain ratios (Twiss and Moores, 1992). So we can use an arbitrary scale or unit for the quadratic elongation ( $\lambda'$ ) axis, as proposed by Ramsay (1967). By determining the position of  $\lambda' = 1$  on the  $\lambda'$  horizontal axis we can obtain a numerical scale for determination of other strain parameters. Construct a Mohr circle with center  $C$  and radius  $R_{XZ}/2$ . The two points where this circle cuts the  $\lambda'$ -axis represent the two principal axes of the strain ellipsoid. At  $\lambda'_3$  construct a line making an angle  $\Phi_{XZ} = 35^\circ$  (Table 1) with the  $\lambda'$ -axis. This line cuts the Mohr circle at the point  $m$ . Drop a perpendicular from  $m$  onto the  $\lambda'$ -axis, intersecting it at  $n$ . This line shows the locus of no finite longitudinal strain (n.f.l.s) and cuts the  $\lambda'$ -axis at  $\lambda' = 1$ . For the plane strain ellipsoid ( $k = 1$ ), the n.f.l.s. locus is a circular section. In all other cases ( $k \neq 1$ ), the two loci do not coincide and the locus of n.f.l.s. represents a surface, not a plane. Then at  $\lambda'_3$  construct a line making an angle  $\Phi_{YZ} = 20^\circ$  with the  $\lambda'$ -axis so it cuts the  $mn$  line (Fig. 3a–c) at  $o$ . The intersection of these lines is located in the YZ principal plane of the strain ellipsoid. Because  $\lambda'_3$  and  $o$  lie on the circumference of  $\lambda'_2 \lambda'_3$  principal plane, the line  $o\lambda'_3$  is a chord of that circle, and its perpendicular bisector intersects the  $\lambda'$ -axis at the center of the  $\lambda'_2 \lambda'_3$  circle in Mohr space. The XY principal plane can be easily determined by constructing a circle with a radius of  $\lambda'_1 \lambda'_2 / 2$ . The magnitude of the principal quadratic elongations  $\lambda'_1$ ,  $\lambda'_2$  and  $\lambda'_3$  can be found by scaling off the distance from the origin; hence:

$$\lambda'_1 = \frac{c\lambda'_1}{cn} \quad \lambda'_2 = \frac{c\lambda'_2}{cn} \quad \lambda'_3 = \frac{c\lambda'_3}{cn}$$

The second case in Table 1 is for a prolate strain ellipsoid. Construct the Mohr circle with center  $C$  and radius  $R_{XZ}/2$  (Fig. 3b). From the minimum principal quadratic elongation ( $\lambda'_3$ ) construct a line making an angle  $\Phi_{XZ} = 30^\circ$  with the  $\lambda'$ -axis. As in the previous case, this line cuts the Mohr circle at  $m$ . Drop a perpendicular from  $m$  onto the  $\lambda'$ -axis, intersecting it at  $n$ . The line  $mn$  cuts the  $\lambda'$ -axis at  $\lambda' = 1$ . Construct an angle  $90 - \Phi_{XY}$  from  $\lambda'_1$  so that it cuts the line of n.f.l.s (or  $mn$ ) at  $o$ . Because it is not possible to find the position of  $\lambda'_2$  on the  $\lambda'$ -axis, it is necessary to use the  $90 - \Phi_{XY}$  angle for constructing  $\lambda'_1 \lambda'_2$  on the circle. The line  $o\lambda'_1$  is a chord of the  $\lambda'_1 \lambda'_2$  circle and its perpendicular bisector intersects the  $\lambda'$ -axis at the center of the  $\lambda'_1 \lambda'_2$  circle ( $q$ ). With this center ( $q$ ) and the radius  $q\lambda'_1$ , one can draw the  $\lambda'_1 \lambda'_2$  circle. Then applying the above equations, it is possible to determine the strain ratios in the XZ, XY and YZ principal planes.

The third case (Table 1) is for a strain ellipsoid with plane strain geometry. All states in this case are the same as that described for



**Fig. 3.** Measuring the strain ratio in  $R_{XY}$  and  $R_{YZ}$  principal planes based on Mohr circle construction with application of the geometry of n.f.l.s and the  $R_{XZ}$  strain ratio for general flattening (a), general constriction (b) and plane strain (c).

constrictional strain and use the  $90-\Phi_{XY}$  angle for constructing the  $\lambda'_1 \lambda'_2$  circle (Fig. 3c). In the cases of uniaxial flattening and constriction, the angle  $\Phi_{YZ} = \Phi_{XZ}$  and  $\Phi_{XY} = \Phi_{XZ}$  and because  $\lambda'_2 = \lambda'_1$  and  $\lambda'_3 = \lambda'_2$ , calculating the strain ratios in the XY- and YZ-planes can be done based on the strain ratio in XZ-plane.

### 3. Practical applications: regional geological background

Deformed conglomerate and metamorphic rocks of the Deh Vazir area in southwestern Iran (Fig. 4a and b) form part of the Sanandaj-Sirjan metamorphic zone (Stöcklin, 1968), within the Zagros orogenic belt of Iran. The Zagros belt, as part of the Himalayan mountain chain, extends for about 2000 km in a NW–SE direction from the East Anatolian fault of Eastern Turkey to the Oman line in southern Iran (Alavi, 1994). The Zagros Orogen formed by continental collision between the Afro-Arabian continent and the Iranian microcontinent in Late Cretaceous to Tertiary time (Berberian and King, 1981). The orogenic belt is the result of closure of Neo-Tethys by consumption of oceanic crust at a NE-dipping subduction zone below the Iranian microcontinent and subsequent Late Cretaceous continental collision

between the Afro-Arabian continent and Iranian microcontinent (Ricou, 1971; Takin, 1972; Dewey et al., 1973; Stocklin, 1968; Berberian and King, 1981; Alavi, 1994; Blanc et al., 2003; McQuarrie, 2004; Sarkarinejad et al., 2008; Sheikholeslami et al., 2008). The Zagros orogenic belt from northeast to southwest consists of three NW–SE trending parallel zones (Fig. 4a): (1) the UrumiehDokhtar Magmatic Belt (UDMB), (2) the Sanandaj-Sirjan HP-LT/HT-LP Metamorphic Belts (SSMB) and (3) the Zagros Fold-and-Thrust Belt (ZFTB). The Sanandaj-Sirjan HP-LT is 150–200 km wide and more than 1500 km long from NW (Sanandaj) to SE (Sirjan) in the western part of Iran (Fig. 4a). Based on metamorphic grade, the Sanandaj-Sirjan zone is subdivided into high-pressure/low temperature and high temperature/low pressure paired metamorphic belt (Sarkarinejad, 1999). The tectonics of the Sanandaj-Sirjan HP-LT metamorphic belt are characterized by numerous thrusts, all transporting rock units from NE to SW in piggyback style (Alavi, 1994).

The results  $^{40}\text{Ar}/^{39}\text{Ar}$  step-heating measurements on biotite, muscovite and amphibole in the Sanandaj-Sirjan HP-LT metamorphic belt are consistent with the overprinting relationship determined from field observations (Sarkarinejad et al., 2009). The first generation of biotite yields plateau ages of  $119.95 \pm 0.88$  and  $112.58 \pm 0.66$  Ma. These late Aptian ages are related to early thrusting and the formation of high-pressure metamorphic rocks at the peak of metamorphism prior to obduction of the Neyriz ophiolite (Sarkarinejad et al., 2009). The Deh Vazir deformed conglomerate is sandwiched between thrust sheets which are part of the Zagros Thrust System (Sarkarinejad and Azizi, 2008; Sarkarinejad et al., 2010), which consists of eight sheets of NW-striking, NE-dipping dextral strike-slip duplex structures that are linked with imbricate fans and oblique slip thrusts (Sarkarinejad and Azizi, 2008).

The most abundant rocks in the study area are deformed conglomerate and micro-conglomerate. The conglomerate pebbles consist of quartzite, phyllite, mica schist; the thickness of this unit varies between 2300 and 2500 m. The matrix around the pebbles is composed of muscovite, quartz, and feldspar. The similarity of composition between pebbles and matrix indicates low rheological contrast between them (Sarkarinejad et al., 2010). The metamorphic grade in this conglomerate is greenschist facies conditions (Sarkarinejad, 1999, 2007).

#### 3.1. Meso- and micro-scale structures

The Deh Vazir conglomerate is strongly foliated and lineated (Fig. 5). The foliation is defined by alignment of pebbles flattened in the XY-plane. The orientation of the foliation varies between  $N80^\circ\text{W}$ ,  $30^\circ\text{NE}$  and  $N30^\circ\text{W}$ ,  $70^\circ\text{NE}$ . The stretching lineation within the foliation has a plunge and trend varying from  $10^\circ$ ,  $N60^\circ\text{W}$  to  $30^\circ$ ,  $N20^\circ\text{W}$ . Stair-stepping structures, mantled  $\sigma$ -type porphyroclast systems with sigmoidal structures, asymmetrical rotated domino boudins, asymmetrical tapering boudins and asymmetrical complex rotated domino and tapering boudins all indicate a top-to-the-SE sense of shear (Fig. 6).

#### 3.2. Structural observations and finite strain analysis

Detailed structural observations and sampling for strain analysis were carried out along a narrow (1 km) section across the thrust faults (Fig. 4). Five structural domains (D1 to D5) with different structural characteristics are defined based on distance from thrust faults (Fig. 7). In each domain an oriented sample was taken for strain analysis and more than 60 foliation and stretching lineation measurements were made for stereographic studies. In strain studies the evidence for volume change is commonly equivocal (Simpson, 1981; Mohanty and Ramsay, 1994), but in some cases, volume change can be discounted (Srivastava et al., 1995;

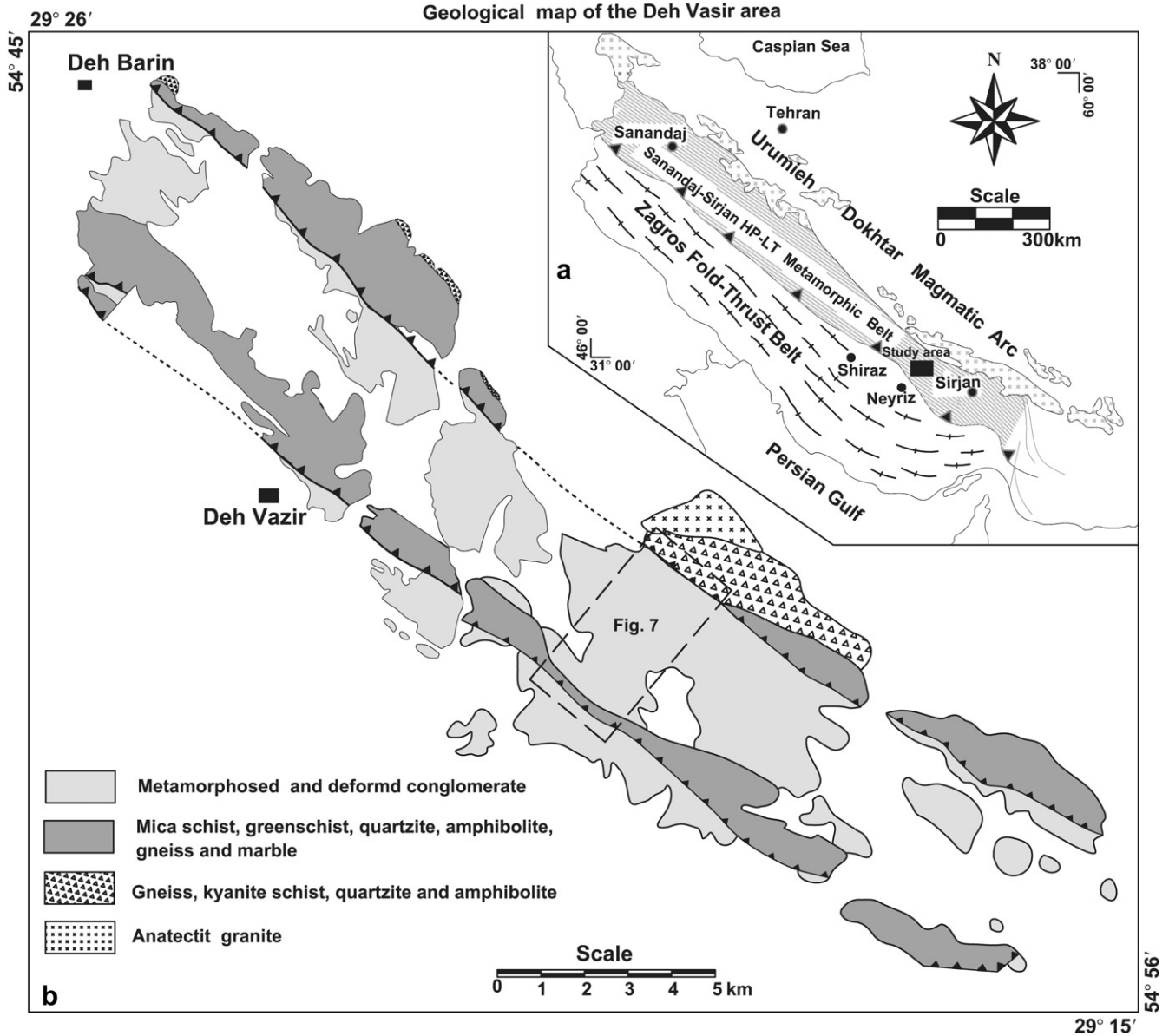


Fig. 4. Geological map of the study area.

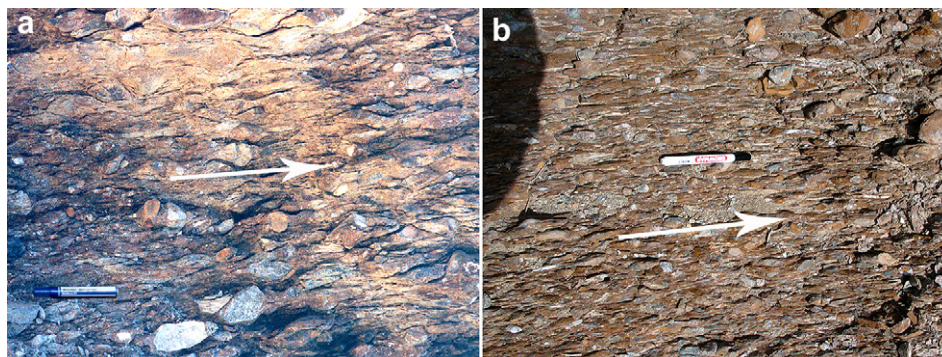
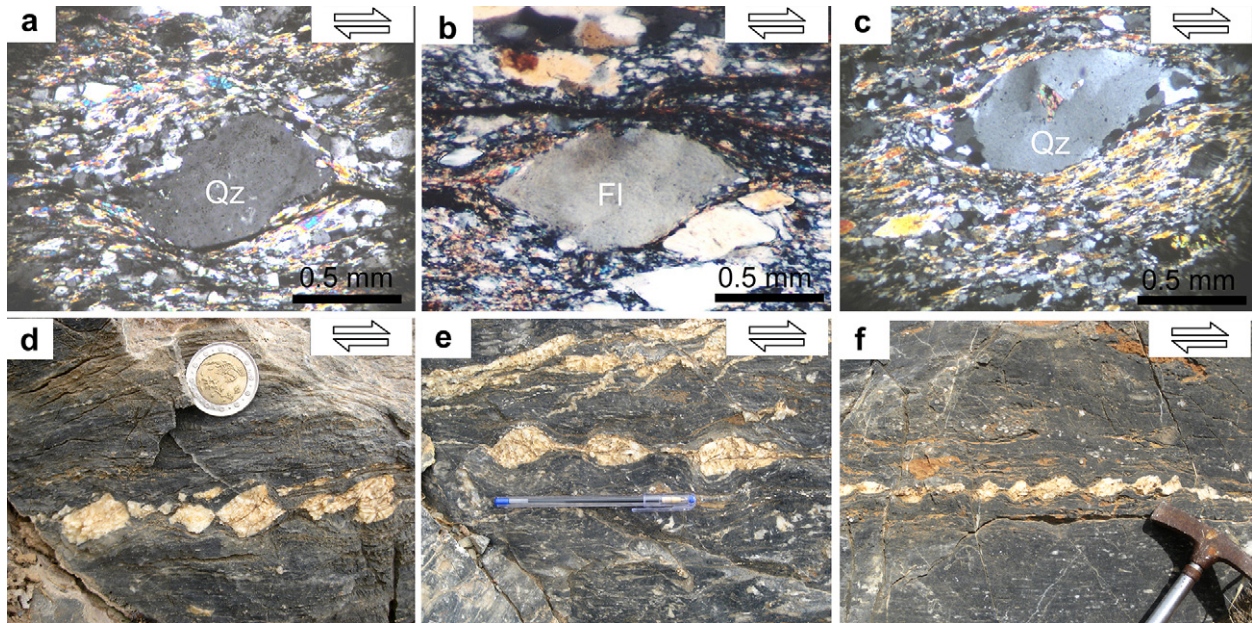


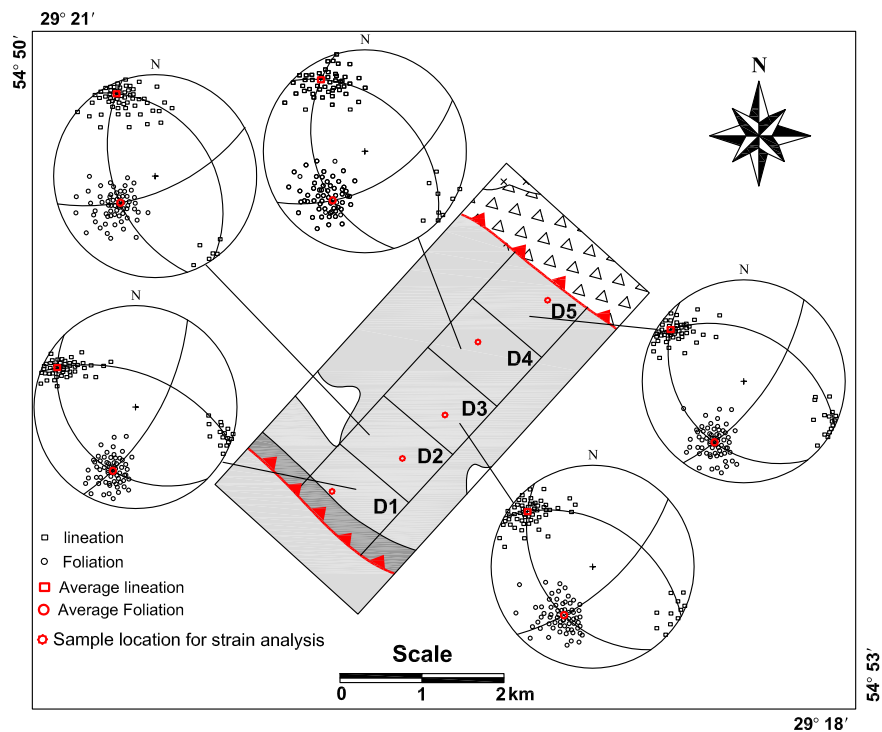
Fig. 5. Preferred orientation of pebble long axes in the Deh Vasir deformed conglomerate. Size of the markers on the left and center of the photos (a and b) is 14 cm. The orientation of the long axes of the pebbles is  $10^\circ, N60^\circ W$  (The photos have been taken looking at inclined faces approximately parallel to the XZ-plane of the finite strain ellipsoid).



**Fig. 6.** Shear sense indicators in the study area. (a, b and c) Photomicrographs showing dextral shear of quartz and feldspar  $\sigma$ -type porphyroclast systems. The matrix consists of muscovite, quartz and feldspar. (d) Dextral asymmetrical rotated domino boudins. (e) Asymmetrical tapering boudins. (f) Asymmetrical complex domino and tapering boudins.

Bhattacharyya and Huddleston, 2001). Volume change during deformation can affect the shape of the finite strain ellipsoid (Ramsay and Wood, 1973). In the Deh Vazir deformed area, the lack of prominent veining implies that volume change was small (Sarkarinejad et al., 2010). In the quartzite pebbles, crystal-plastic deformation was dominant as indicated by subgrain-rotation recrystallization microstructures (Sarkarinejad and Azizi, 2008) and there is little or no microstructural evidence for solution

transfer, also indicating approximately constant volume deformation. Previous strain studies in this area (Sarkarinejad, 2007; Sarkarinejad and Azizi, 2008) on the deformed conglomerates and micro fossils of the Sanandaj-Sirjan HP-LT metamorphic belt approximately show the plane strain geometry of finite strain ellipsoid ( $K \approx 0.85$ ) with only small volume changes. Quartz c-axis fabrics are type-I crossed-girdle patterns, which indicate approximately plane strain ( $K = 1$ ) conditions (Sarkarinejad and Azizi,



**Fig. 7.** Five structural domains (D1 to D5), sample locations and distribution of foliation and lineation on lower hemisphere equal area stereographic projections using SpheriStat 2.2 software.

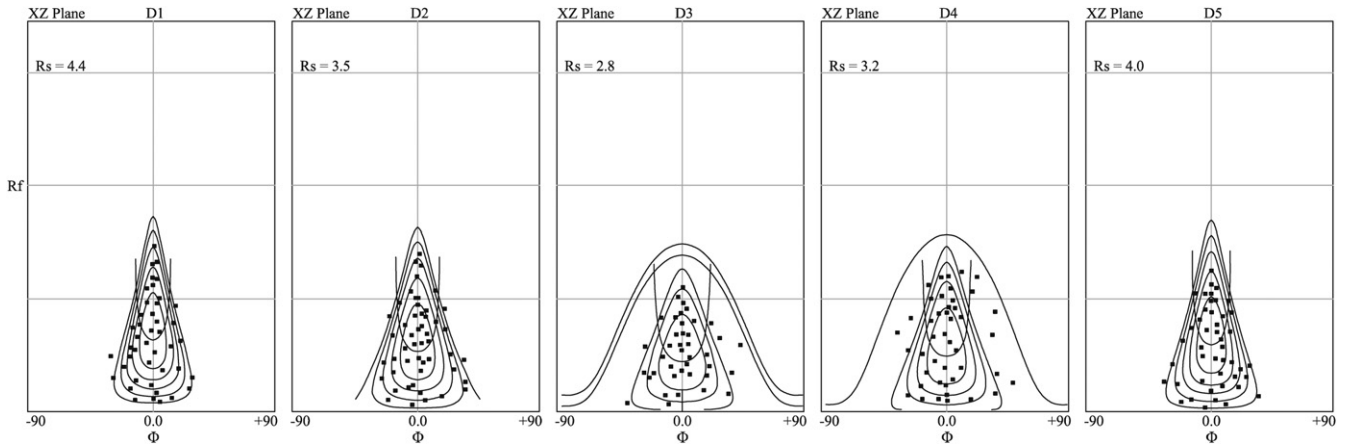


Fig. 8.  $Rf/\phi$  diagrams in the XZ-plane for B1 to B5 samples.

2008; Sarkarinejad et al., 2010). Therefore throughout this analysis, we have assumed constant volume deformation.

As mentioned by Xypolias (2009), the kinematic vorticity parameter ( $W_m$ ) plays an important role in determining the orientation of the principal axes of the finite strain ellipsoid with respect to the instantaneous stretching axes (ISA). Kinematic vorticity is a dimensionless measure of rotation relative to strain and

characterizes the amount of shortening relative to displacement.  $W_m$  was originally defined as an instantaneous rotation relative to the instantaneous stretching at a point (Truesdell, 1953; Means et al., 1980). Most of the vorticity methods utilize data collected on the XZ-plane of finite strain (parallel to lineation and normal to foliation) and commonly assume steady-state deformation with the vorticity vector approximately parallel to the Y-axis of the strain ellipsoid.

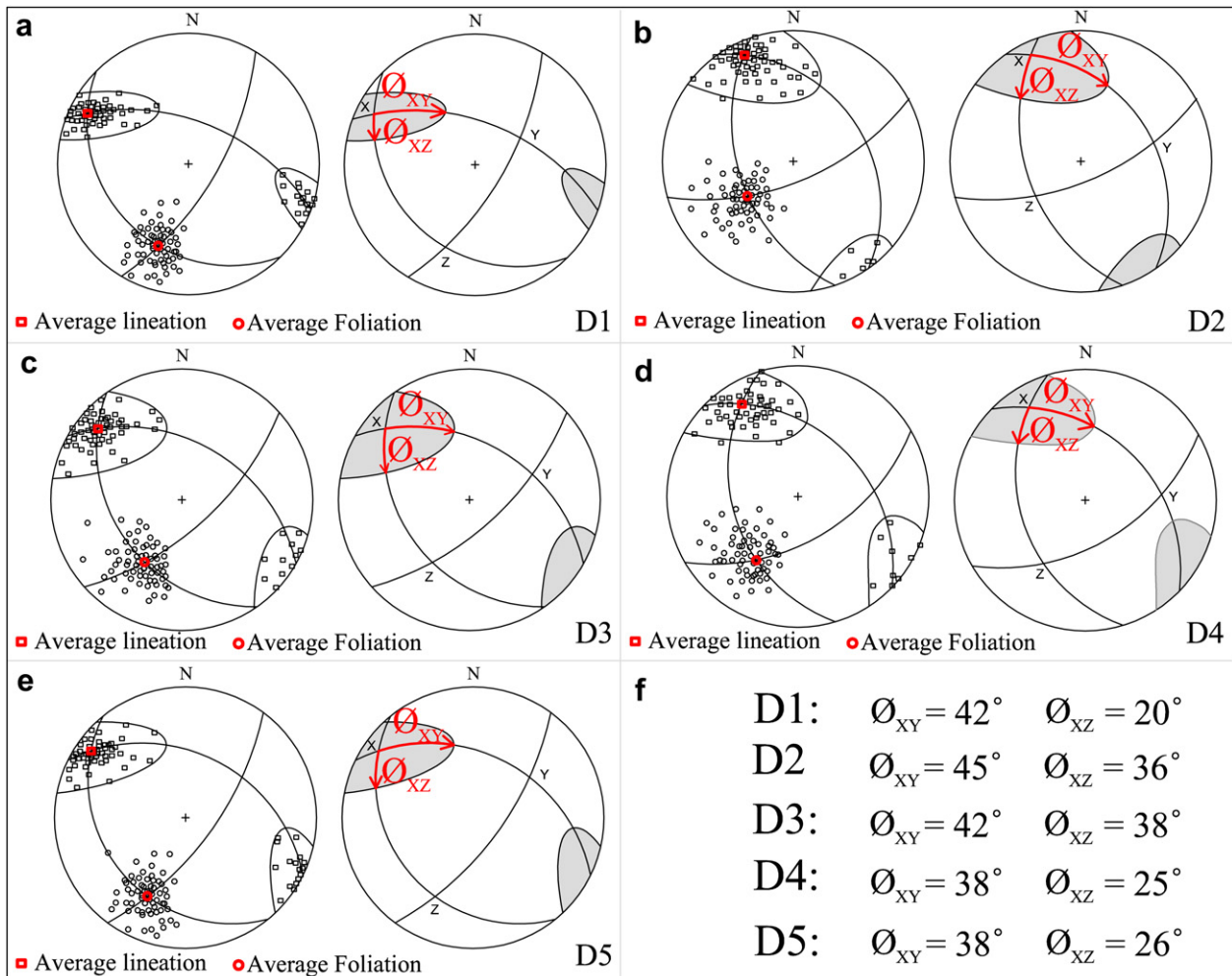


Fig. 9. Lower hemisphere, equal area projection of the n.f.l.s geometry ( $\phi_{XY}$  and  $\phi_{XZ}$ ) based on foliation and stretching lineation structural elements (The stereonets have been plotted using SpheriStat 2.2 software).

**Table 2**  
Strain ratios on the XY and YZ principal planes of the strain ellipsoid.

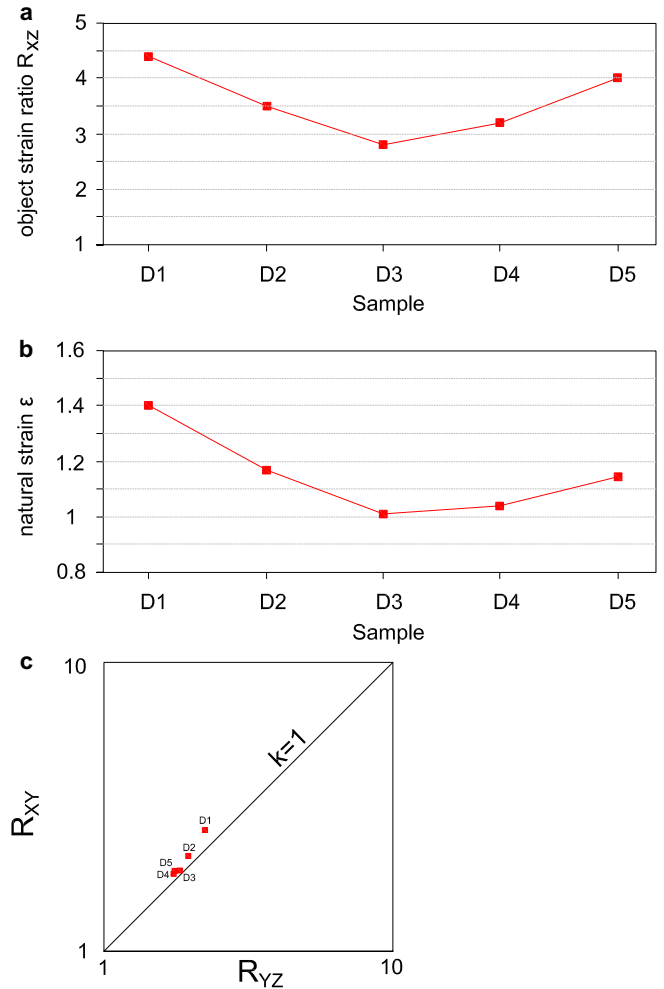
	$R_{XZ}$	$\Phi_{XZ}$	$\Phi_{XY}$	$R_{XY}$	$R_{YZ}$
D1	4.4	20°	42°	4.2	3.5
D2	3.5	36°	45°	3.3	2.9
D3	2.8	38°	42°	2.8	2.6
D4	3.2	25°	38°	2.7	2.4
D5	4.0	26°	38°	2.8	2.5

For cases of simple shear and sub-simple shear,  $W_m$  is measured on a scale between 0 and 1, with 0 being pure shear and 1 being simple shear. The  $W_m$  scale is not linear, but can be converted to a linear scale by considering the percent of a deformation resulting from simple shear and pure shear. Fort and Bailey (2007) propose three separate fields for pure, general, and simple shear dominated deformations. Pure shear dominated deformations have  $W_m$ -values of 0–0.3, corresponding to less than 20% simple shear. In contrast, simple shear dominated deformations have  $W_m$ -values of greater than 0.95, corresponding to greater than 80% simple shear. General shear occupies the range between 0.3 and 0.95.

$W_m$  is an important factor indicating fit or discord of principal axes of the finite strain ellipsoid and the instantaneous stretching axes of deformation. In pure shear dominated deformation, the maximum and minimum instantaneous stretching axes ( $ISA_1$  and  $ISA_2$ ) approximately coincide with the short and long axes of the finite strain ellipse (Passchier and Trouw, 2005; Xypolias, 2009) and the deformation is mainly coaxial. In this case the pole of foliation and orientation the stretching lineation show the directions of shortening and extension, respectively. Kinematic vorticity analysis of the Deh Vazir deformed area revealed a prominent pure shear component of deformation (Sarkarinejad et al., 2010; Samani, 2010). So in this study we assumed that the shortening and extension directions approximately coincide with the pole to foliation and the stretching lineation, respectively. Moreover in this framework we can use other structures such as fold axial planes and boudin necks for separation of shortening and extension domains (Talbot, 1970).

Strain analysis was performed on a representative sample from the deformed conglomerate in each of the 5 domains (D1 to D5); sample localities are shown on Fig. 7. All samples were collected from conglomerate layers containing low competency contrast between pebbles and matrix and well defined planar and linear fabric elements. The pebble shape was used as the strain marker. Length-to-width ratios of pebbles were determined from measurements made on sections cut normal to the foliation and parallel to the lineation.

The study relies on the following assumptions: (1) the foliation plane coincides with the XY-plane of the strain ellipsoid, (2) the stretching lineation defines the long axis of the strain ellipsoid, (3) the strike of the main thrust faults defines the direction of the shear zone boundaries and  $\phi$  is the angle between the X-axis of the elliptical strain marker and the reference line of the shear zone boundary trace and (4) the deformation was isochoric (constant volume). The strain ratios ( $R_{XZ}$ ) were estimated for each sample applying the  $R_f/\phi$  method (Ramsay, 1967; Lisle, 1985) (Fig. 8). In order to determine the 3D geometry of the finite strain ellipsoid, the above described method was used to estimate the tectonic strain ratios in the XY and YZ principal planes. By applying stereographic analysis of the stretching lineation and foliation, we determined the geometry of the surfaces of n.l.f.s. More than 60 orientations of foliation and stretching lineation were measured at each sample location in the D1 to D5 domains. For each domain the distributions of orientations were plotted on lower hemisphere, equal area stereographic projections (Fig. 7). The orientations of n.l.f.s. were determined. Their geometry can be completely



**Fig. 10.** (a and b) Variation of finite strain parameters ( $R_{XZ}$  and  $\epsilon$ ) plotted against sample locations from the thrust planes. (c) Ellipsoid shape analyzed by plotting the finite strains for XY and YZ principal sections on a Ramsay diagram.

described by the angles made by their lines of intersection on the principal planes and one of the principal strain axes (angles  $\Phi_{XY}$ ,  $\Phi_{XZ}$ ,  $\Phi_{YZ}$ ) (Fig. 9). Finally by using the geometry of n.l.f.s. (angles  $\Phi_{XY}$ ,  $\Phi_{XZ}$ ,  $\Phi_{YZ}$ ) and the tectonic strain ratio on the XZ-plane and by applying the method described above, the strain ratios on the XY- and YZ-planes were calculated (Table 2). These three values of strain (for the XZ-, XY- and YZ-planes) have been used to evaluate the finite natural logarithmic strain (Ramsay and Huber, 1983):

$$\epsilon = \left(\frac{1}{3}\right)^{1/2} \left[ (\ln(R_{XZ}))^2 + (\ln(R_{YZ}))^2 + (\ln(R_{XY}))^2 \right]^{1/2}$$

In order to investigate the finite strain variation across the shear zones, finite strain parameters ( $R_{XZ}$  and  $\epsilon$ ) were plotted against sample locations from the thrust planes (Fig. 10a and b); for each distribution, curves of mean values have been plotted. The ellipsoid shape was analyzed by plotting the finite strain for the XY- and YZ-planes on a Ramsay diagram (Fig. 10c).

#### 4. Discussion

The shear zones associated with the Zagros thrust system (Sarkarinejad and Azizi, 2008) are characterized by a top-to-the-SW sense of shear and lengthening parallel to the shear direction. This conclusion is based on two main lines of evidence: (1) the



stretching lineation is roughly parallel to the shear direction as obtained by independent field evidence (Fig. 6); and (2) the stretching lineation trend is at a high angle with respect to the general dip direction of the foliation (Fig. 9). *XZ* finite strain values obtained from deformed pebbles, as well as the calculated natural logarithmic strain ( $\epsilon$ ), indicate that finite strain increases in an approximately linear fashion toward the thrust faults (Fig. 10a and b). Similar linear trends have been observed in other shear zones (Talbot and Sokoutis, 1995). Vitale and Mazzoli (2008) discriminated mylonite types using strain intervals of  $\epsilon = 0-1$  (proto-mylonites),  $\epsilon = 1-2.5$  (mylonite) and  $\epsilon > 2.5$  (ultramylonite). Our analyzed samples have strain values ( $1 < \epsilon < 1.4$ ) corresponding to mylonite. The Ramsay diagram (Fig. 10c) shows that most of the object finite strain ellipsoids inferred from the deformed pebbles fall into the prolate field, although they plot close to the plane strain line. The finite strain ellipsoid distribution in the Ramsay diagram is similar to that obtained for the region by applying  $R_{XZ}$ ,  $R_{XY}$  and  $R_{YZ}$  strain ratios in planes cut parallel to the finite strain principal planes (Sarkarinejad et al., 2010). Thus both previous methods and the newly described method of finite strain determination give approximately similar values for the study area.

## 5. Conclusion

3D strain analysis based on the geometry of the surface of no finite longitudinal strain and the strain ratio measured on the *XZ*-plane is a useful method for strain studies in deformed areas with approximately homogeneous deformation. This method is especially useful for situations when preparation of 3 perpendicular cut sections is difficult or impossible.

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